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# Effect of mechanical strain on lateral photovoltaic effect in n-3C-SiC/n-Si hetorjunction Toward mechanical strain sensors capable of photoenergy harvesting

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#### ABSTRACT

It is beneficial to have sensors capable of harvesting photoenergy in the era of Internet of Things (IoT) and 5 G integrated infrastructure. In this paper, we report the piezo-optoelectronic coupling characteristic of n-type 3C-SiC/n-type Si heterostructure and demonstrate it with a mechanical strain sensing device capable of harvesting photoenergy. The device is capable of generating a photovoltage as high as 3.4 mV when illuminated by a laser power as small as 10  $\mu$ W, i.e. four times higher than similar sensors in previous studies. In addition, the strain sensitivity  $\Delta V/V/\epsilon$  ratios are 9.94 and 13.11 for tensile and compressive strain, respectively, which are five times better than conventional metal strain gauges. These findings contribute to the overall understanding of the piezo-optoelectronic coupling effect of SiC/Si heterojunction paving the way for the development of multifunctional sensors with light harvesting capabilities.

#### 1. Introduction

In the era of advanced infrastructure driven by the integration of Internet-of-Things (IoT) and 5 G technologies, it is critical to monitor the surrounding environment through an extensive network of sensors which provide consistent real-time sensing data, allowing us to make proper predictions or decisions that contribute to maintaining a safe, secure and convenient living environment [1]. However, a conventional physical-wired sensor system requires multiple wires to power and transmit sensing data. This approach is costly and labour-consuming to install and maintain, and for some remote area, this becomes a near-impossible challenge. Therefore, there is a high interest in sensors that can harvest energy from surrounding environment and transmit sensing information wirelessly [2,3]. A practical method for sensors to harvest the surrounding photoenergy is by utilizing lateral photovoltaic effect (LPE), which was first discovered by Schottky in 1930 [4] and later developed by Wallmark in 1957 [5]. Since then, the LPE effect has been extensively studied in various semiconductor materials for energy harvesting application [6–14].

Among the materials presenting LPE effect, silicon carbide (SiC) is a wide bandgap material exhibiting many attractive properties such as superior mechanical, chemical, optical and electrical performance [15]. The wide bandgap property not only helps SiC become advantageous in applications requiring high power [16], high working temperature [17], and high frequency [18], but also enables the development of SiC sensors capable of measuring signals in conditions where conventional Si sensors show degradation, such as high-temperature sensor [19], accelerometer [20], or photodetector [12,21–23]. However, the costly fabrication process and small size of 4H-SiC and 6H-SiC present significant challenges, hindering SiC development for applications requiring many sensors [24].

Coupling between multiple physical effects is considered a potential method to enhance the performance of the available sensing techniques. One example is the coupling among piezoelectricity and optoelectronic, known as the piezo-phototronic effect, which enhances the performance of optoelectronic devices by introducing strain modulation in ZnO material [25]. Another method is coupling between piezoresistivity and optoelectronic, known as piezo-optoelectronic effect, has attracted

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interest for further enhancing the piezoresistive effect by modulating photogenerated charge carriers in SiC/Si heterojunction [26]. Table S1 in the Supporting Information compares between these coupling methods and traditional mechanical strain sensing techniques.

Recent studies on the piezo-optoelectronic effect focus mainly on ptype SiC/Si heterojunction, as p-type SiC shows better piezoresistive performance than n-type SiC material [27]. However, this leaves a knowledge gap that requires investigation into the mechanical, electrical, and optoelectronic properties of n-type SiC/ n-type Si heterojunction. Additionally, n-type SiC/ n-type Si heterojunctions possess many beneficial properties that should not be overlooked. Firstly, n-type SiC is easier to fabricate using unintentionally doped films method, whereas p-type SiC fabrication requires high temperature (approximately 1350°C), which is nearly the melting temperature of Si [28]. Secondly, in terms of electron-hole pairs separation effect, which is an important effect governing the performance of LPE, the heterostructure of n-type SiC/ n-type Si has more advantages than n-type SiC/ p-type Si heterostructure. Considering n-SiC is the top layer of heterojunction, n-Si substrate is more beneficial for the separation of electrons-holes at the heterojunction than p-Si substrate. Particularly, in the case of n-Si substrate, when laser light is absorbed in Si substrate, more electrons contribute to the generation of LPE, as the generated electrons have a lower recombination probability due to less holes available in the n-Si substrate. In the contrary, for p-Si substrate, the generated electrons easily recombine with the holes available within the p-Si substrate. Therefore, using n-Si substrate produces a better LPE as more electrons are separated at the heterojunction. Because of the above reasons, we decided to investigate and report on the mechanical and electrical properties of n-type SiC/ n-type Si heterojunction.

In this paper, we report the piezo-optoelectronic coupling characteristic of n-type SiC/n-type Si heterostructure and demonstrate it with a mechanical strain sensing device capable of harvesting photoenergy. The ability to collect photoenergy of the device is characterized through lateral photovoltage output measurement, and then the impact of coupling mechanical strain on lateral photovoltaic effect is investigated. The experimental data reveals that the sensing structure exhibits superior light harvesting capabilities, generating a lateral photovoltage as high as 14 mV over 40  $\mu$ W of laser power, which is three to four times larger than that observed in previous studies, along with excellent mechanical strain sensing capabilities, presenting a  $\Delta V/V/\epsilon$  ratios of 9.94 and 13.11 when applying tensile strain and compressive strain, respectively, which is five times higher than that of conventional metal strain gauges. The underlying mechanism for the piezo-optoelectronic coupling effect is explained based on the photon absorption and carrier diffusion theory. Therefore, the investigation results will significantly contribute to the overall understanding of the properties of SiC material, especially n-type SiC/n-type Si heterojunction, and open up new possibilities for the fabrication of multifunctional sensors with light harvesting capabilities.

The research results and discussions are structured as follows. First, the sensors design concepts, fabrication process, and experimental setup are explained. Then, the photoenergy harvesting capability is characterized. Next, the impact of mechanical strain on the lateral photovoltaic effect is investigated, followed by discussion of the underlying physical explanation. Finally, the highlights of the research are mentioned in the conclusion.

# 2. The concept of strain sensing using piezo-optoelectronic effect

Fig. 1 illustrates the concept for the light harvesting strain sensor constructed from n-type SiC thin layer grown on n-type Si substrate, forming an n-3C-SiC/n-Si heterostructure, with two aluminium electrodes positioned on the top surface of the SiC thin film. The SiC layer in SiC/Si heterojunction is beneficial for the light harvesting performance, as the SiC layer allows harvesting a wider wavelength spectrum compared to Si/Si homojunction. Additionally, by combining the wide bandgap SiC with narrow bandgap Si, the heterostructure exhibits a large built-in potential which in turn improves the separation of EHPs at the SiC/Si heterojunction [29].

An external laser source of wavelength 637 nm is used to illuminate one end of the sensing element. The SiC layer, with its wide bandgap property, is mostly transparent to the laser wavelength 637 nm, allowing laser light to penetrate into the Si substrate below and be absorbed by the Si substrate [30]. This addition photoenergy leads to the generation of electron-hole pairs (EHPs) within the Si substrate. The generated EHPs are then separated due to the built-in voltage at the SiC/Si heterojunction, resulting in electrons drifting upward to SiC layer and holes drifting downward to Si substrate.

At the illuminated area, the electron concentration within the SiC layer increases (point A), while those at the non-illuminated area remains low (point B). This result in excess electrons diffusing laterally from point A to point B, with lower electron concentration. Two aluminium electrodes positioned at point A and B collect the electrons at these locations. The difference in electrons concentration between the two points results in a lateral photovoltage,  $V_{AB}$ , measured between the two electrodes (Fig. 1a).

Applying tensile or compressive strain to the sensing area between the two electrodes alters the electron mobility within the SiC layer, consequently affecting the electron concentration gradient between A and B (Fig. 1b,c). We can estimate the level of strain applied to the sensing area by measuring the changes in the output voltage  $V_{AB}$  before and after strain application. This sensing element can be used to constantly monitor the amount of strain applied to the sensor and simultaneously harvest the photoenergy from the external laser source illuminated onto the device.



Fig. 1. : Concept of the light-harvesting strain sensor using piezo-optoelectronic effect in case of a) no strain applied, b) under tensile strain and c) under compressive strain.

#### 3. Fabrication and experimental setup

We have fabricated a proof-of-concept device, and the fabrication process comprises eight steps as illustrated in supplementary Figure S1. First, a commercialize n-type Si wafer with a thickness of 380 µm, phosphorous doped with doping concentration of  $1 \times 10^{15}$  atom/cm<sup>3</sup> was used. The Si wafer was thoroughly cleaned following the standard Radio Corporation of America (RCA) cleaning procedure. Then, singlecrystalline unintentionally n-type doped 3C-SiC thin films were epitaxially grown on the clean Si wafer using the low-pressure chemical vapor deposition process (LPCVD) at 1000°C, employing propylene (C3H6) and silane (SiH<sub>4</sub>) as precursors [28]. The thickness of SiC layer was measured using Nanometrics Nanospec 210 instrument, and the value was approximately 500 nm. The doping concentration of SiC layer was measured using hot probe technique [31] and the value was approximately  $3 \times 10^{17}$  atom/cm<sup>3</sup> (Step 2). In the next step, aluminium (Al) was sputtered on top of SiC layer to a thickness of 500 nm, following by a photoresist layer deposited on the top using spin coating method. The electrode designs were subsequently patterned to the photoresist layer by exposing it to ultraviolet light using maskless aligner MLA150 of Heidelberg Instruments. The exposed area became soluble when submerged into developer solution (Step 3–5). We used wet etching method to remove the Al which was not covered by the photoresist mask, forming two Al electrodes on top of the SiC layer with dimensions 200 µm width, 500 µm length, and separated 500 µm from each other (Step 6). The piezoresistive effect in SiC is anisotropic; therefore, to achieve the highest sensitivity, the aluminium electrodes were aligned in <100> orientation, where largest piezoresistive performance is observed for n-type SiC [32]. The excess photoresist was removed using Tegal 915 and the wafer was diced into individual cantilever devices of 40 mm length and 9 mm width (Step 7). Al thin wires were used in the wire bonding process to connect between the electrodes on the device with an external electrical printed circuit board (PCB), improving the ability to handle the device during the characterization experiments (Step 8).

The schematic and geometry of the device are illustrated in supplementary Figure S2a. Supplementary Figure S2b shows a Transmission Electron Microscopy (TEM) image of a cross-section of the device, showing the microstructure at the interface of SiC/Si heterojunction and the thickness of the grown epitaxy 3C-SiC layer was approximately 500 nm. The selected area diffraction (SAED) patterns in supplementary Figure S2c confirm the single crystallinity of the grown 3C-SiC layer.

To characterize the piezo-optoelectronic effect of the heterostructure, we used the bending cantilever method, and the experimental setup is illustrated in Fig. 2. One end of the cantilever was fixed by a cshape clamp, while the other free-end was hung downward or pull upward using weights, introducing tensile strain or compressive strain to the sensing area, respectively. The weights used were 20 g, 40 g and 60 g, corresponding to the strain values of  $154 \times 10^{-6}$ ,  $308 \times 10^{-6}$  and  $462 \times 10^{-6}$ , respectively. The electrical properties were characterized using Keithley 2450 source meter.

The photo-response was tested using laser with wavelength of

637 nm. The laser was mounted on an XYZ stage (PT3, Thorlabs) to fix the laser spot's position inside a u-shape cut out on the clamp. This setup ensured that the laser position remained stable during the strain application experiments. The laser spot diameter was measured using BC100006N-VIS Beam profiler (Thorlabs) and the spot diameter value was 100  $\mu$ m. Laser power was monitored using S130C power sensors (Thorlabs) and PM100D power meter (Thorlabs). All experiments were conducted under dark conditions and at room temperature (25°C).

#### 4. Result and discussion

First, the current-voltage (I-V) characteristics of the device under dark condition were measured and then the light harvesting ability was evaluated under laser illumination at three powers:  $20 \ \mu$ W,  $40 \ \mu$ W, and  $60 \ \mu$ W (Fig. 3a). The linearity of all four I-V curves indicates that contacts between two aluminium electrodes and the SiC layer are ohmic contacts. Under dark conditions, the I-V curve went through the origin of the coordinates. electrode A of the device with a red laser (wavelength 637 nm), the I-V curve was offset from the origin due to the generation of lateral photovoltage/photocurrent [33].

Fig. 3b shows the generation of the lateral photovoltage (LPV) output under non-uniform illumination of red laser at three powers 20 µW, 40  $\mu$ W, and 60  $\mu$ W. In the case of no laser illumination (dark condition), the lateral photovoltage output was 0 mV. Under non-uniform illumination of red laser power 20 µW, the output photovoltage value increased to 7 mV, and kept increasing linearly, reaching 14 mV and 21 mV when increasing the laser power to 40  $\mu$ W and 60  $\mu$ W. We could clearly see the excellent repeatability and stable response of the photovoltage output reflecting the laser status when turned on and off. When we shorted the external circuit, we could measure a photocurrent generated across the sensing element, and the short-circuit photocurrent values are shown in Fig. 3c. Under dark conditions, the generated photocurrent was 0 µA, which increased to 2.3  $\mu$ A under illumination of laser power 20  $\mu$ W. The negative sign shows that the photogenerated electrons moved within SiC layer from electrode A to B. The generated photocurrent further increased to 4.6  $\mu$ A and 6.9  $\mu$ A when increasing the laser power to 40  $\mu$ W and 60  $\mu$ W. In Fig. 3d, it can be seen clearly that the output photovoltage scaled linearly with laser power from 0  $\mu$ W to 500  $\mu$ W, and eventually reached saturation photovoltage of 97 mV at laser power around 600 μW. Further increasing laser power beyond this point yielded negligible improvement in the photovoltage output. Compared with the previous studies, the lateral photovoltage achieved in this study is approximately three to four times larger than that reported in previous studies [12,33]. The lateral photovoltage performance comparison with other similar devices is summarised in Supplementary Table S2.

The underlying mechanism for the lateral photovoltage output is illustrated in Fig. 4, with Fig. 4a depicting the SiC/Si heterojunction in dark condition, and Fig. 4b-d illustrate the lateral photovoltage forming mechanism when non-uniformly illuminating a red laser beam to the device.

Fig. 4a shows the device consisting of two layers: an n-type SiC layer



Fig. 2. The experiment setup using bending cantilever method to characterize the piezo-optoelectronic coupling effect of the SiC/Si device.



Fig. 3. a) Current-voltage (I-V) characteristic of the device. b) Repeatability test of lateral photovoltage (LPV) generation under illumination of red laser with different power and c) Repeatability test of photocurrent generation under illumination of different laser power. d) Generation of LPV under illumination of laser power increasing from 10  $\mu$ W to 2000  $\mu$ W.

deposited on an n-type Si substate. The difference in electron concentration between the Si and SiC layer leads to the movement of electrons and holes across the interface of the two materials [12,34]. Electrons from the SiC layer diffuse through the interface to the Si substrate, leaving positive charges in the SiC layer. Conversely, holes from Si substrate diffuse to the SiC layer, resulting in negative charges in the Si substrate. This migration of electrons and holes eventually reaches equilibrium, creating the depletion zone, at the interface between the two materials. An electric field ( $E_0$ ) is formed across the depletion zone, preventing further diffusion of electrons and holes through the junction.

Fig. 4b illustrates the underlying mechanism of LPV output generation under non-uniform illumination of red laser wavelength 637 nm to the device. When the red laser light wavelength 637 nm illuminates to the device, it initially contacts with SiC layer. SiC is a wide bandgap material with a bandgap value around 2.38 eV, which is larger than the photoenergy of red laser, value of approximately 1.95 eV. Due to its wide bandgap property, the SiC layer is mostly transparent to the red laser, resulting in the laser passing through the SiC layer and penetrating into the Si substrate below to a depth of approximately 4.2  $\mu$ m [35] from the SiC/Si interface. The bandgap of Si is 1.12 eV, which is suitable for the absorption of red laser, resulting red laser photons being absorbed by the Si substrate [30]. The laser photons excite the electrons in the Si substrate, allowing them to jump to the conduction band, leaving holes

back in the valence band, resulting in generation of electron-hole pairs (EHPs) in Si substrate under the laser spot.

Fig. 4d illustrates the energy band diagrams of the heterojunction at the cross section of electrodes A and B. The built-in electric field across the junction separates the generated electron-hole pairs, causing electrons to drift up toward the SiC layer (opposite to the direction of the built-in electric field), while pushing the holes to drift down towards the Si substrate.

This results in an increase of electron concentration in SiC layer and holes in Si substrate at the laser point, electrode A. However, at nonilluminated area (electrode B), the electron concentration in SiC layer remain unchanged as there are minimal additional EHPs generated. According to the absorption theory [36,37], when illuminate light of frequency  $\nu$  to a semiconductor with energy gap E<sub>g</sub>, the light-excited electrons concentration at the illumination point can be written as:

$$n_0 = K_1 (h\nu - E_g)^{\alpha} \tag{1}$$

where  $K_1$  is a proportional coefficient, h is the Planck constant,  $\alpha$  is an exponential coefficient. Some of these excited electrons have a chance, R, to recombine with the holes, or re-excited  $\tau p/n(0)$  times in average, where  $\tau$  is the lifetime of diffusion electrons. The electron concentration within SiC layer at laser point is defined as [38]:



**Fig. 4.** a) The mechanism of SiC/Si heterojunction form within the device in dark condition and b) The mechanism of lateral photovoltage generation under nonuniform illumination of laser. The LPE includes two phenomena: the generation of electron holes pairs (EHPs) under laser illumination and the electrons diffusion from point A to B within SiC thin layer. c) Diffusion of electrons from A to B due to difference in electron concentration between the two points d) The energy band diagram of SiC and Si layers at laser illuminated region (electrode A), and non-illuminated region (electrode B).

$$N_0 = n_0 [1 - R^{(\tau p/n(0))}]$$
<sup>(2)</sup>

Subsequently, the excess electrons in the SiC layer diffuse laterally from the illuminated region (electrode A) to the surrounding circular areas (electrode B) due to the differences in the electron concentration. Along the diffusion path, some electrons interact with holes and recombine, while the remaining electrons are collected at the electrodes A and B. The electrons diffuse laterally within SiC layer to the surrounding area could be described with diffusion model [30,39]:

$$D\frac{d^2N_x}{dx^2} = \frac{N_x}{\tau_{SiC}}$$
(3)

where  $\tau_{SiC}$  is the lifetime of diffusion electrons in SiC. The electron concentration at a certain point distance *x* on the diffusing path can be defined by:

$$N_x = N_0 \times \exp\left(-\frac{x}{d}\right) \tag{4}$$

where  $N_0$  is the electron concentration at the illumination point, *x* is the distance from the laser point to the probing position, and  $d = (D_{SiC} \times \tau_{SiC})^{1/2}$  is the diffusion length of electron in SiC, with  $D_{SiC}$  is the diffusion constant of SiC.

The electron diffusion profile is qualitatively described in Fig. 4c. At the laser illumination point (electrode A), the distance from laser spot to point A is  $x_A = 0$ , result in the electron concentration at point A is  $N_A = N_0$ . When *x* value moves further away from the illumination point (electrode B), the distance from laser spot to point B is  $x_B > 0$ , the

exponential  $e^{\left(-\frac{x}{d}\right)}$  value quickly decrease, leading to the exponential decay of the electron concentration. Therefore, the electron concentration at point B becomes significantly smaller compared to point A. shows the difference in electron concentration between point A and B result in the voltage V<sub>AB</sub> measured across the electrodes, and can be calculated as [39]:

$$V_{AB} = K_2 N_0 \left[ \exp\left(-\frac{|L-x|}{d}\right) - \exp\left(-\frac{|L+x|}{d}\right) \right]$$
$$\approx \frac{2K_2 N_0}{d} \times \exp\left(-\frac{L}{d}\right) x$$
(5)

where  $K_2 = \frac{k_B T}{n_0}$  is the proportional coefficient, *x* is the laser position, and 2*L* is the distance between the two electrodes.

The mechanical sensing performance of the device was further characterized using weights hanging at the free-end of the cantilever. Force was applied to the free-end of the cantilever using three different weights: 20 g, 40 g and 60 g, resulting in tensile strain ( $\varepsilon$ ) values of  $154 \times 10^{-6}$ ,  $308 \times 10^{-6}$  and  $462 \times 10^{-6}$  exerting at the sensing area. Without laser illumination (i.e. in dark condition) these strain values could not be detected, as there was no lateral photovoltage across the sensing area. However, the strain sensing element is activated when non-uniformly illuminating the clamped-end of the cantilever with laser, allowing the lateral photovoltage generation and simultaneously measuring mechanical strain applied to the sensing area.

The strain values applied to the sensor are measured based on the difference between the photovoltage outputs before and after the strain applied. Consider the photovoltage output in the case of free strain as reference, the output photovoltage difference ( $\Delta V$ ) is shown in Fig. 5.

At laser power 20  $\mu$ W, the output photovoltage in the case of freestrain was 7 mV (Fig. 3b). Introducing tensile strain to the sensing area reduced the photovoltage output by  $-11 \mu$ V,  $-22 \mu$ V, and  $-32 \mu$ V, corresponding to the three weight values 20 g, 40 g and 60 g, respectively (Fig. 5a). Using higher laser power resulted in the photovoltage reducing more significantly (Fig. 5b,c).

Fig. 5d shows an excellent linear relation ( $R^2 \approx 99.76\%$ ) between the output photovoltage and tensile strain ( $\epsilon$ ) varying from 0 to  $500 \times 10^{-6}$ , with the sensitivity of the mechanical strain sensing element was 0.07  $\mu$ V/ $\epsilon$ , 0.137  $\mu$ V/ $\epsilon$  and 0.207  $\mu$ V/ $\epsilon$  corresponding to laser powers of 20  $\mu$ W, 40  $\mu$ W and 60  $\mu$ W, respectively. It is worth noting that for all three laser power settings 20  $\mu$ W, 40  $\mu$ W and 60  $\mu$ W, we observed a consistent result in fractional changes ( $\Delta$ V/V) of lateral photovoltage output. The absolute of  $\Delta$ V/V increased linearly with the applied strain, with the ratios of  $\Delta$ V/V/ $\epsilon$  was 9.94 when tensile strain applied. (Fig. 7a).

Subsequently, a pulley system was used to change the force direction applied to the free end of the cantilever. Three different weights: 20 g, 40 g and 60 g were used to pull up the free end of the cantilever, corresponding to compressive strain ( $\varepsilon$ ) values of  $154 \times 10^{-6}$ ,  $308 \times 10^{-6}$  and  $462 \times 10^{-6}$  at the sensing area, respectively. Consider the photovoltage output at free strain as reference, the output photovoltage difference ( $\Delta V$ ) is shown in Fig. 6. Fig. 6a-c clearly show that introducing compressive strain to the sensing area increased the photovoltage output, and the mechanical strain sensitivity also increased when using higher laser power. The sensitivity of the mechanical strain sensing element was 0.088  $\mu V/\varepsilon$ , 0.181  $\mu V/\varepsilon$  and 0.267  $\mu V/\varepsilon$  according with the laser power 20  $\mu W$ , 40  $\mu W$  and 60  $\mu W$ , respectively (Fig. 6d).

Fig. 7 shows the linear relationship between fractional changes of

photovoltage output and the tensile strain or compressive strain value varies from 0 to  $500 \times 10^{-6}$  at three different laser powers 20  $\mu$ W, 40  $\mu$ W and 60  $\mu$ W. With all three laser power settings, we can easily see the excellent linear relationship of fractional changes ( $\Delta$ V/V) of photovoltage output across the range of applied strain. The ratios of  $\Delta$ V/V/ $\epsilon$  was 9.94 for tensile strain and 13.11 for compressive strain, respectively, demonstrating excellent mechanical strain sensing sensitivity which is five times higher than conventional metal strain gauges, which have a typical ratio value only approximately 2 [40].

Based on the theory of energy valleys shifting and electron transfer in SiC conduction band [41], we can qualitatively explain the mechanism for the piezo-optoelectronic coupling effect and illustrated in Fig. 8. Under free-strain condition, the energy valleys have equal energy level and the output voltage measured across electrodes A and B is mainly due to the lateral photovoltage effect (Fig. 8a).

Under a tensile strain along [100] orientation, the energy valleys in direction [100] (i.e., the longitudinal valleys) shifts upward, whereas the energy valleys in direction [010] and [001] (i.e., the transverse valleys) shift downward [41]. These shifts in the six energy valleys result in the repopulation of free electrons from the valleys with high energy to lower energy valleys, meaning that the electrons from the longitudinal direction are redistributed to the valleys in transverse direction. Moreover, since the mobility of electron in transverse direction ( $\mu_{\perp}$ ) is higher than that in the longitudinal direction ( $\mu_{\parallel}$ ) [32], more electrons in the transverse direction ([010] and [001]) lead to a higher total effective mobility, meaning that more electrons move faster towards electrode B. Consequently, the resistance between the two electrodes is reduced, resulting in the reduction of the output voltage V<sub>AB</sub> (Fig. 8b). Applying higher tensile strain result in more electrons repopulating in the



**Fig. 5.** (a - c) Repeatability test of the photovoltage output under three different tensile strain ( $\varepsilon$ ) value  $154 \times 10^{-6}$ ,  $308 \times 10^{-6}$  and  $462 \times 10^{-6}$  with three laser powers 20  $\mu$ W, 40  $\mu$ W, and 60  $\mu$ W. d) The linear change of photovoltage output difference ( $\Delta$ V) corresponding to tensile strain ( $\varepsilon$ ) varies from 0 to  $500 \times 10^{-6}$  at three different laser power 20  $\mu$ W, 40  $\mu$ W and 60  $\mu$ W.



**Fig. 6.** (a - c) Repeatability test of the photovoltage output under three different compressive strain ( $\varepsilon$ ) value  $154 \times 10^{-6}$ ,  $308 \times 10^{-6}$  and  $462 \times 10^{-6}$  with three laser powers 20  $\mu$ W, 40  $\mu$ W, and 60  $\mu$ W. d) The linear change of photovoltage output difference ( $\Delta$ V) corresponding with compressive strain ( $\varepsilon$ ) varies from 0 to  $500 \times 10^{-6}$  at three different laser power 20  $\mu$ W, 40  $\mu$ W and 60  $\mu$ W.



Fig. 7. The linear relationship between fractional changes of lateral photovoltage output and applied strain value in the case of a) tensile strain and b) compressive strain using three laser powers 20 µW, 40 µW and 60 µW.

transverse electron, leading to larger output voltage reduction, and therefore larger  $\Delta V$ .

On the contrary, under compressive strain, the energy valleys in direction [100] shift downward, while the energy valleys in direction [010] and [001] shift upward. These shifts result in the repopulation of free electrons from the valleys in transverse direction to the valleys in the longitudinal direction, leading to lower total effective mobility, meaning fewer electrons arrive at electrode B. Consequently, the resistance between the two electrodes is increased, resulting in an increase of

the output voltage V<sub>AB</sub> (Fig. 8c). This hypothesis agrees well with the experimental results where V<sub>AB</sub> decreased ( $\Delta V$  became negative) when tensile strain was applied, and V<sub>AB</sub> increased ( $\Delta V$  became positive) when compressive strain was applied to the cantilever device.

#### 5. Conclusion

In conclusion, we have investigated the impact of mechanical strain on lateral photovoltaic effect in n-type SiC/n-type Si heterojunction and

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Fig. 8. The mechanism of the piezo-optoelectronic coupling effect under a) non-strain, b) tensile strain and c) compressive strain conditions.

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its application for mechanical strain sensors with photoenergy harvesting capabilities. A proof-of-concept cantilever device has been successfully designed and fabricated, featuring a simple vet robust design which consistently sensed the external strain applied while simultaneously harvested the photoenergy exposed to the sensing element. The sensing structure demonstrated superior light harvesting capabilities, generating three to four times larger lateral photovoltage output than that observed in previous studies. Under non-uniform illumination of red laser wavelength 637 nm with power as low as 10  $\mu$ W, the sensing structure was capable of outputting a lateral photovoltage up to 3.4 mV, which kept increasing linearly, reaching 7 mV, 14 mV, and 21 mV when increasing the laser power to 20  $\mu W,$  40  $\mu W$  and 60  $\mu W,$  respectively. The device also demonstrated excellent mechanical strain sensing capabilities, with a strain sensitivity  $\Delta V/V/\epsilon$  ratios was 9.94 for tensile strain and 13.11 for compressive strain. This is five times more sensitive than that of conventional metal strain gauges.

The proposed device exhibits excellent repeatability in light harvesting and strain sensing performance. It also shows great long-term performance stability, as consistent results were observed during the experiments period, and repeatable results were achieved months later. The mechanism for light harvesting and strain sensing capabilities have been explained based on absorption theory and diffusion theory, which consistently aligned with our experimental results.

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Investigating a full theoretical analysis model to quantitatively explain the effect of strain on the energy band structure of SiC under photo excitation is essential for developing future mechanical strain sensors and will be our future works. The findings in this study remarkably contribute to the overall understanding of the piezooptoelectronic coupling effect of SiC/Si heterojunction, paving the way toward the development of multifunctional sensors with light harvesting capabilities.

#### CRediT authorship contribution statement

**Dang D.H. Tran:** Writing – review & editing, Writing – original draft, Visualization, Validation, Methodology, Investigation, Formal analysis, Data curation, Conceptualization. **Tuan-Hung Nguyen:** Writing – review & editing, Resources, Investigation, Formal analysis, Conceptualization. **Dzung Viet Dao:** Writing – review & editing, Validation, Supervision, Resources, Project administration, Funding acquisition. **Cong Thanh Nguyen:** Writing – review & editing, Validation, Resources, Formal analysis. **Erik W. Streed:** Resources, Project

administration, Funding acquisition. **Nam-Trung Nguyen:** Project administration, Funding acquisition. **Van Thanh Dau:** Writing – review & editing, Supervision, Resources, Funding acquisition.

#### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### Data Availability

Data will be made available on request.

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#### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.nanoen.2024.109844.

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